

The consortium



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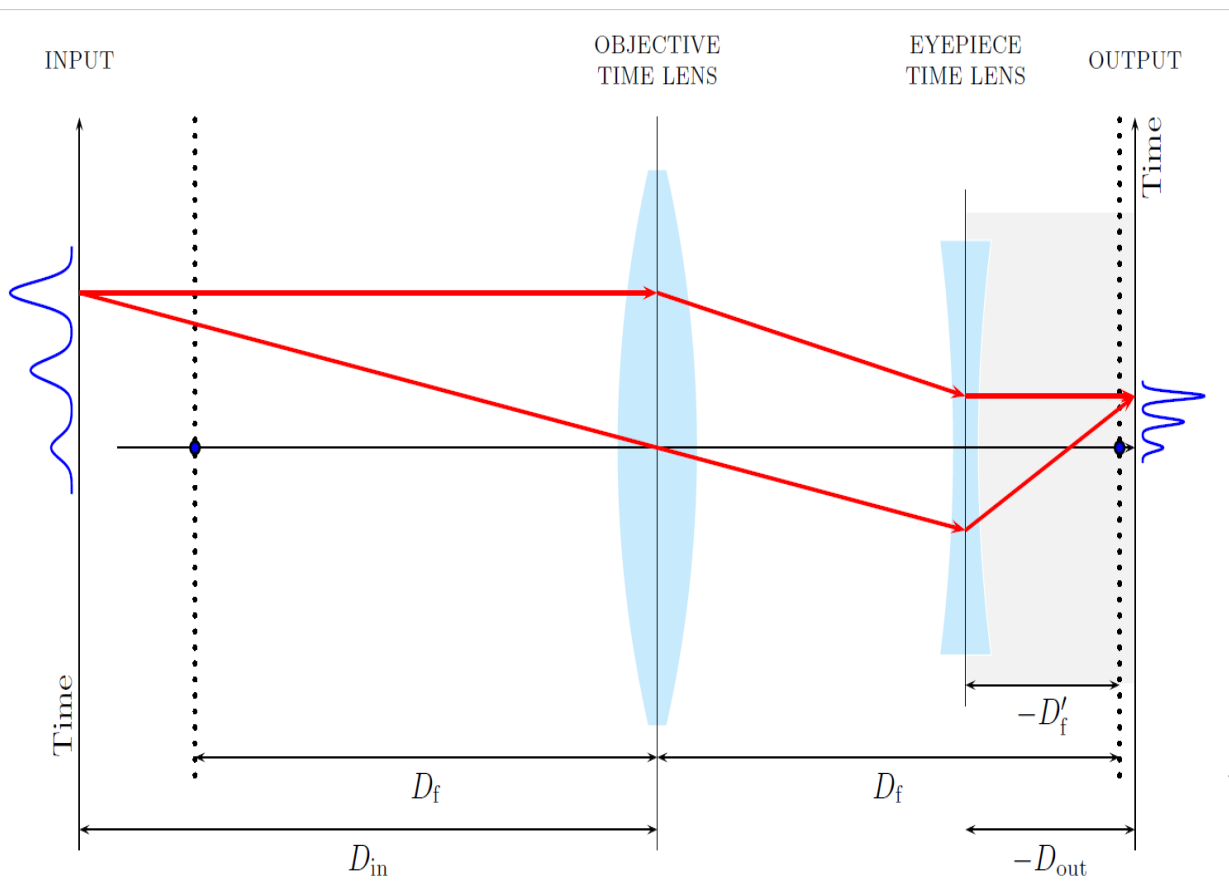


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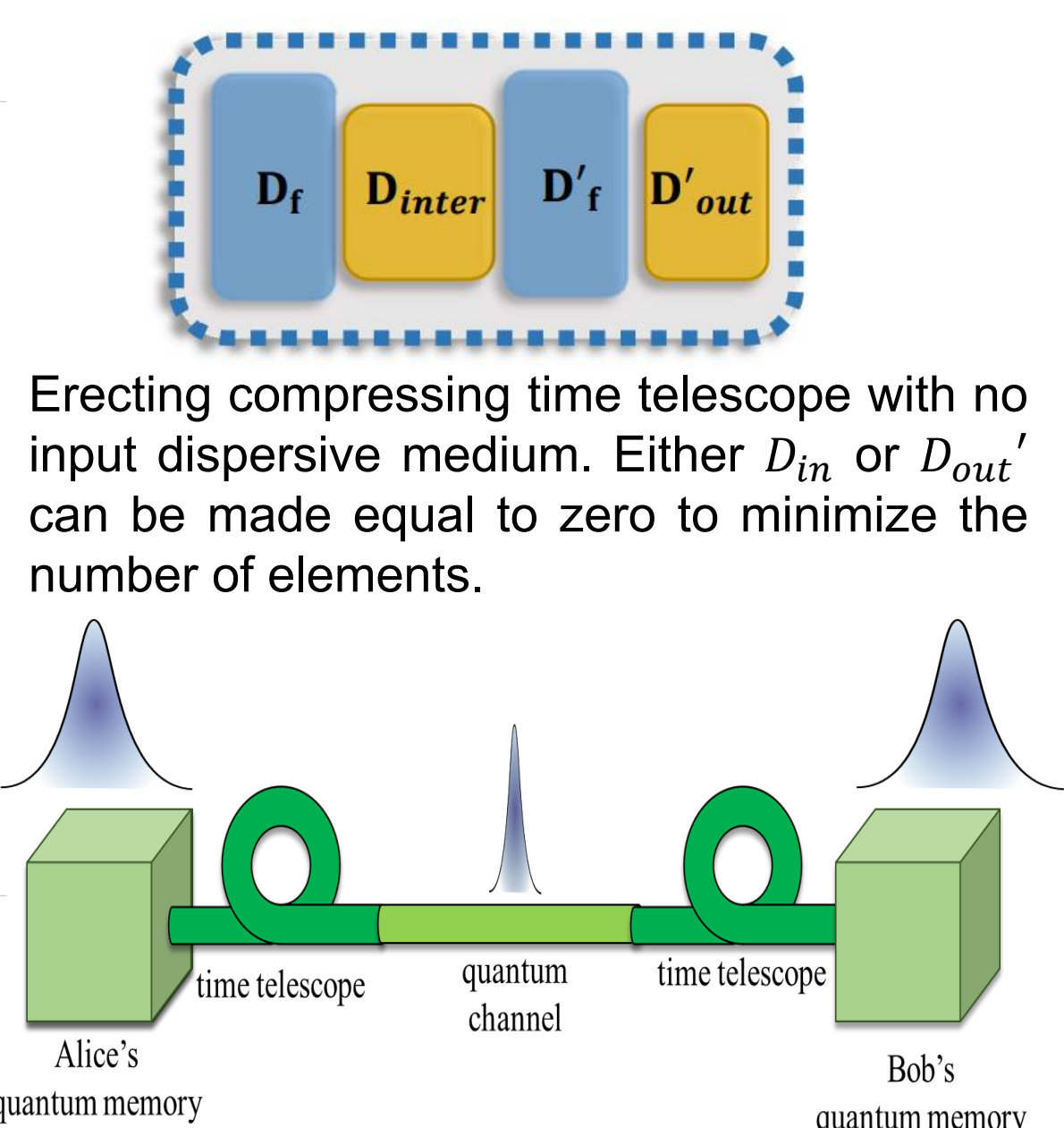
University of Warsaw: M. Karpiński, S. Kapoor, M. Ogrodnik

Quantum temporal imaging: an erecting time telescope



Geometrical optics representation of the erecting compressing time telescope. The grey area shows a dispersive medium with a negative GDD, resulting in the creation of a real image at the output. This element is not possible in the spatial domain because negative diffraction does not exist while negative dispersion does.

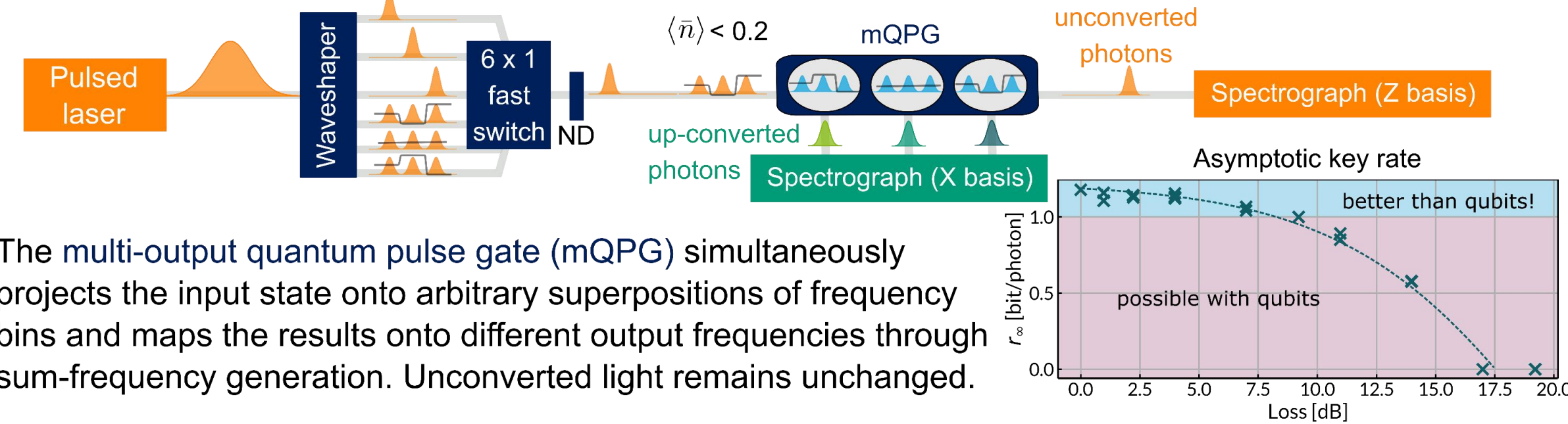
S. Srivastava, D. B. Horoshko, M. I. Kolobov, Opt. Express 31, 38560-38577 (2023)
Recent experimental realization at Bar Ilan Univ.: O. Raphaely et al., Opt. Laser Technol. 192, 113912 (2025).



Erecting compressing time telescope with no input dispersive medium. Either D_{in} or D_{out}' can be made equal to zero to minimize the number of elements.

Converting ps-scale pulses in the telecom band, optimal for high-rate fiber transmission, to nanosecond scale pulses in the visible range processed by quantum memories. The pulses can be made identical leaving the encoded quantum information untouched.

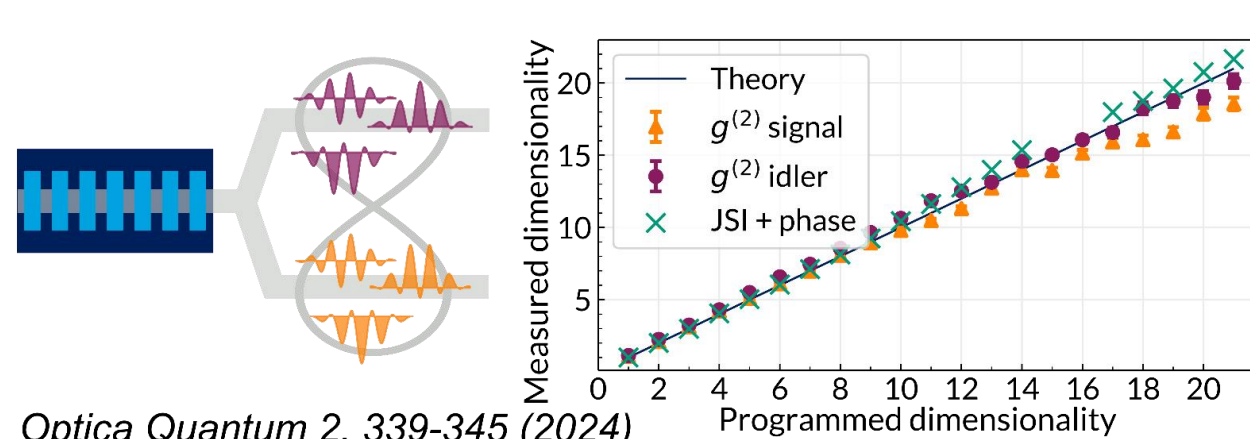
High-dimensional quantum key distribution with temporal modes



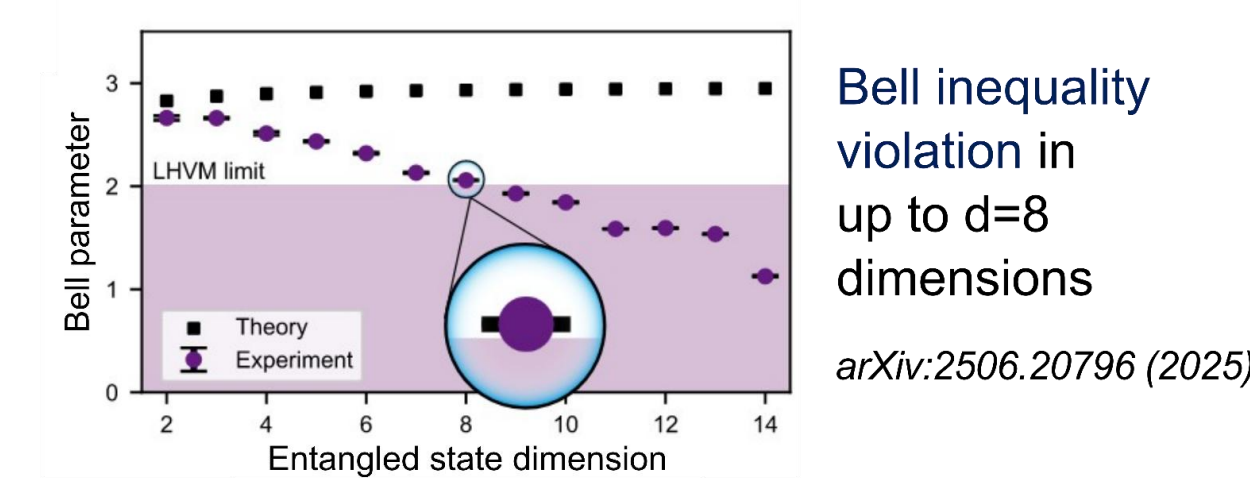
The multi-output quantum pulse gate (mQPG) simultaneously projects the input state onto arbitrary superpositions of frequency bins and maps the results onto different output frequencies through sum-frequency generation. Unconverted light remains unchanged.

High-dimensional entanglement

Programmable source of photons in controlled high-dimensional entangled temporal modes



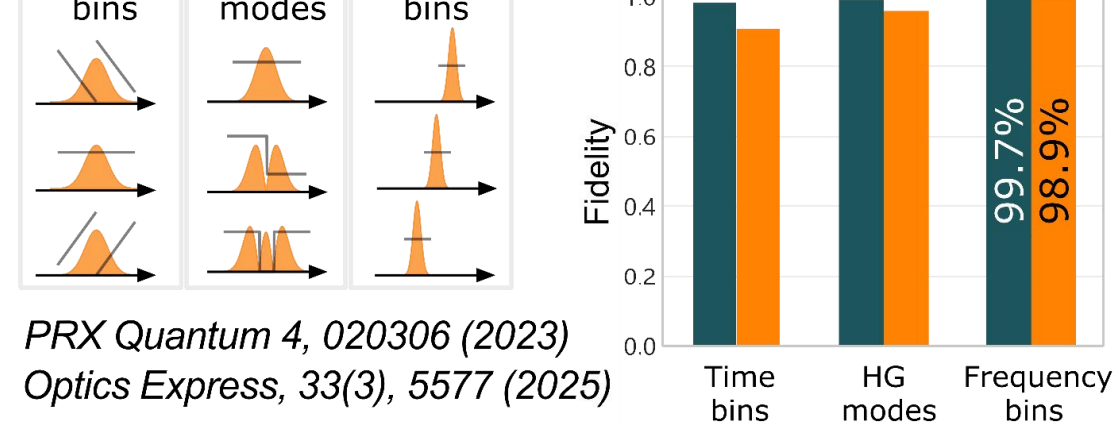
Optica Quantum 2, 339-345 (2024)



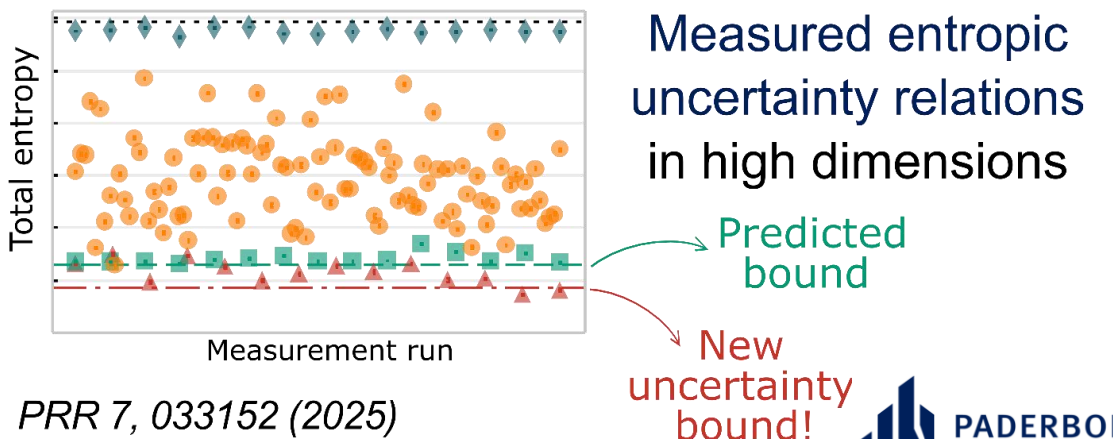
arXiv:2506.20796 (2025)

Programmable mode-sorting

High-fidelity projections in different encodings



PRX Quantum 4, 020306 (2023)
Optics Express, 33(3), 5577 (2025)



PRR 7, 033152 (2025)

Construction of efficient Schmidt number witnesses

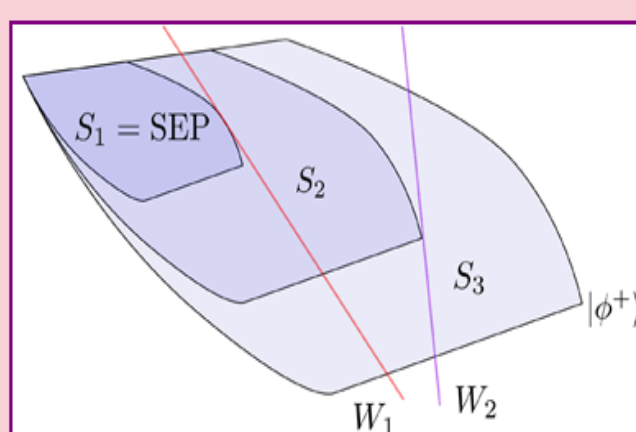


Motivation

- Modern setups (Quantum phase gate, time-frequency encoding, ...) prepare high-dimensional quantum states.
- Problem: **Certify dimensionality** of prepared states.

Schmidt number

- Decompose pure state: $|\psi\rangle = \sum_{i=0}^{k-1} \sqrt{\lambda_i} |i\rangle_A |i\rangle_B$
- k is called **Schmidt rank**, measures dimension of state.
- For mixed states: convex roof yields Schmidt number
- Detect using Schmidt number witness W :
 $\text{Tr}(W\rho_k) \geq 0 \quad \forall \rho \in S_k$,
 $\text{Tr}(W\rho) < 0$ for some ρ .
- Simple witness:
 $W = 1 - \frac{d}{k} |\phi^+\rangle\langle\phi^+|$
with $|\phi^+\rangle = \frac{1}{\sqrt{d}} \sum_{i=0}^{d-1} |ii\rangle$



The algorithm

- Problem: requires $O(d^2)$ measurements in standard basis.
 - Construct witness that uses measurements M with $|M| =$
- Choose k , subset M and some constant $C \in (-1, 1)$.
 - Run semidefinite program
find $\tilde{W}^{(1)}$
s.t. $\langle \phi^+ | \tilde{W}^{(1)} | \phi^+ \rangle = -1$,
 $-1 \leq \tilde{W}^{(1)} \leq 1$.
 - Find $\min_{\phi_k^{(1)}} \langle \phi_k^{(1)} | \tilde{W}^{(1)} | \phi_k^{(1)} \rangle$.
 - Go back to 2., add constraint
 $\langle \phi_k^{(1)} | \tilde{W}^{(2)} | \phi_k^{(1)} \rangle \geq C$.
 - Stop if infeasible (decrease C) or converged (increase C) and repeat.

Result

- Measuring $\text{Tr}(\tilde{W}\rho) < -\sqrt{\frac{d^2 - 4d + 4k}{d^2}}$ with
 $\tilde{W} = \left(1 - \frac{2}{d}\right) |00\rangle\langle 00| - \frac{2}{d} \sum_{i=1}^{d-1} (|00\rangle\langle ii| + |ii\rangle\langle 00|) - \left(1 - \frac{2}{d}\right) \sum_{i=1}^{d-1} |ii\rangle\langle ii|$
- certifies Schmidt rank $k+1$.

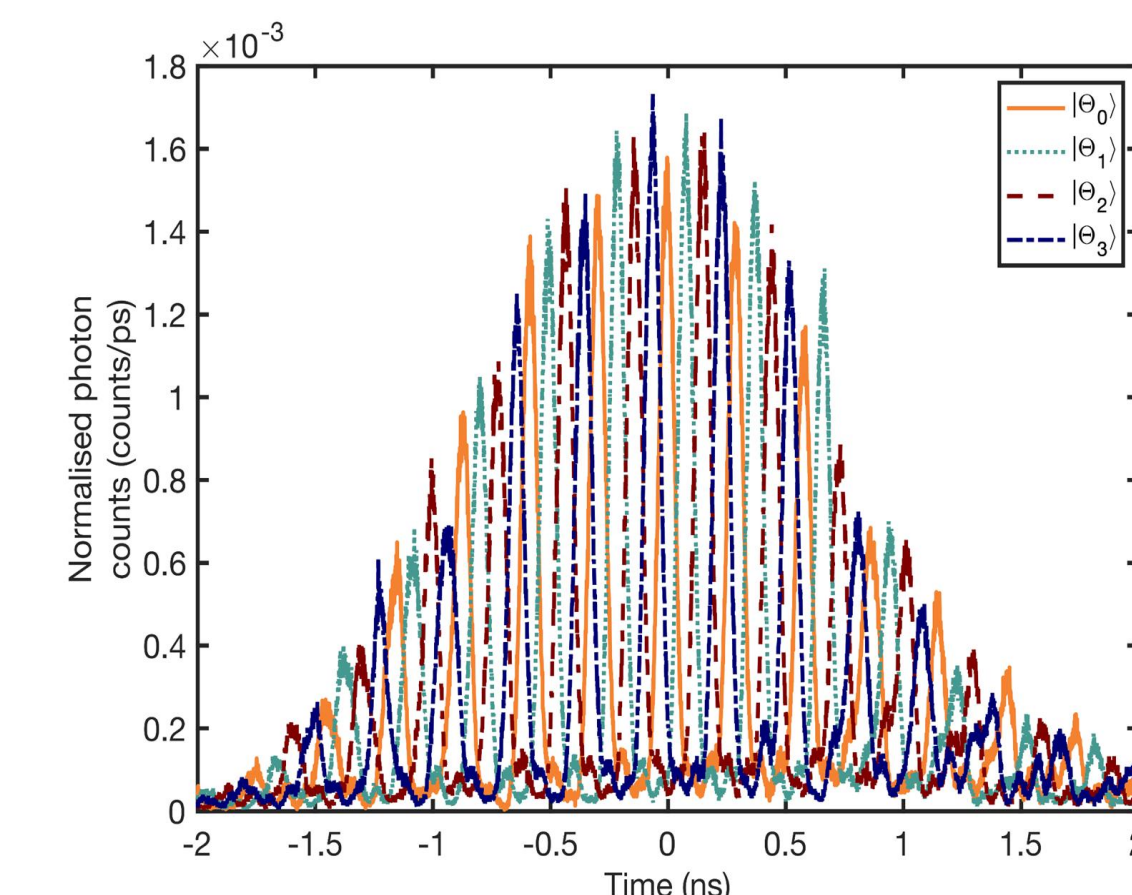
N. Wyderka, G. Chesi, H. Kampermann, C. Macchiavello, D. Bruß, Phys. Rev. A **107**, 022431 (2023)

Efficient detection of multidimensional single-photon time-bin superpositions



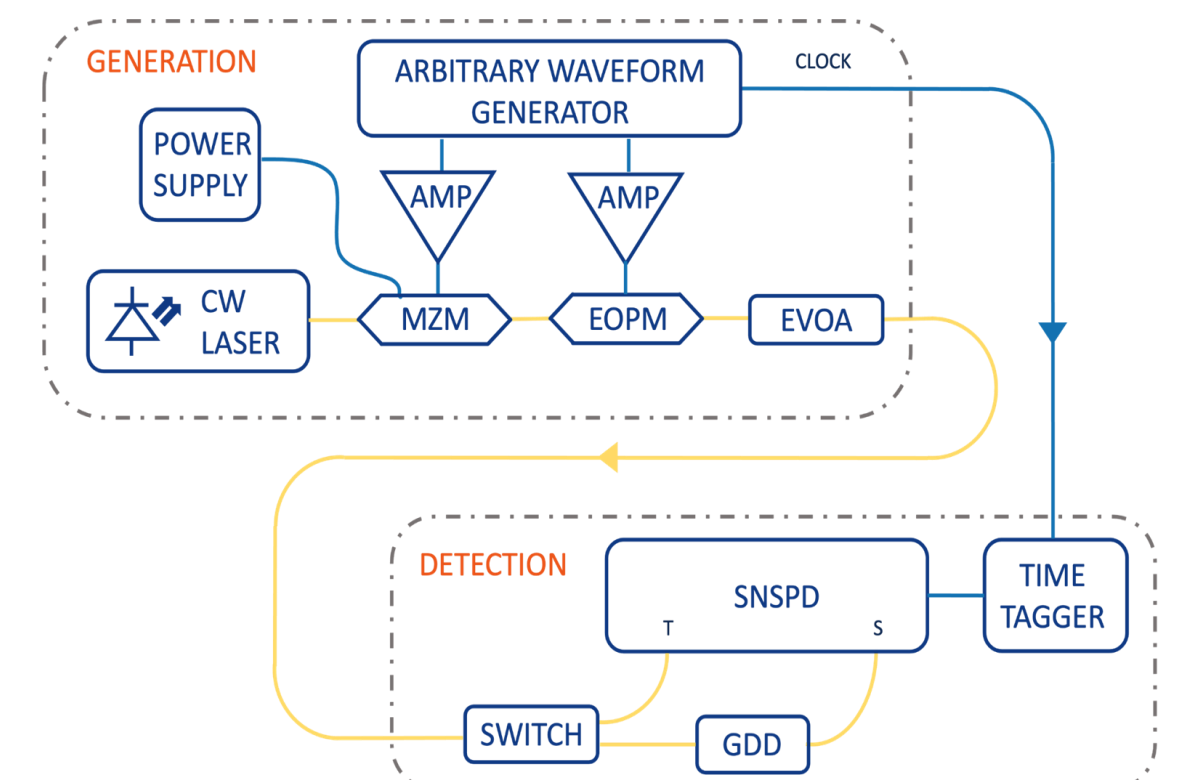
photon.fuw.edu.pl

With dispersive medium we can detect time-bin superpositions in the single-photon-counting regime in an all-fiber setup without the use of interferometers. We showed that we can do that efficiently thanks to the temporal Talbot effect. Currently we are working on using this method for high-dimensional quantum key distribution.



Measured time of arrival histograms for superpositions of four pulses with equal time separation.

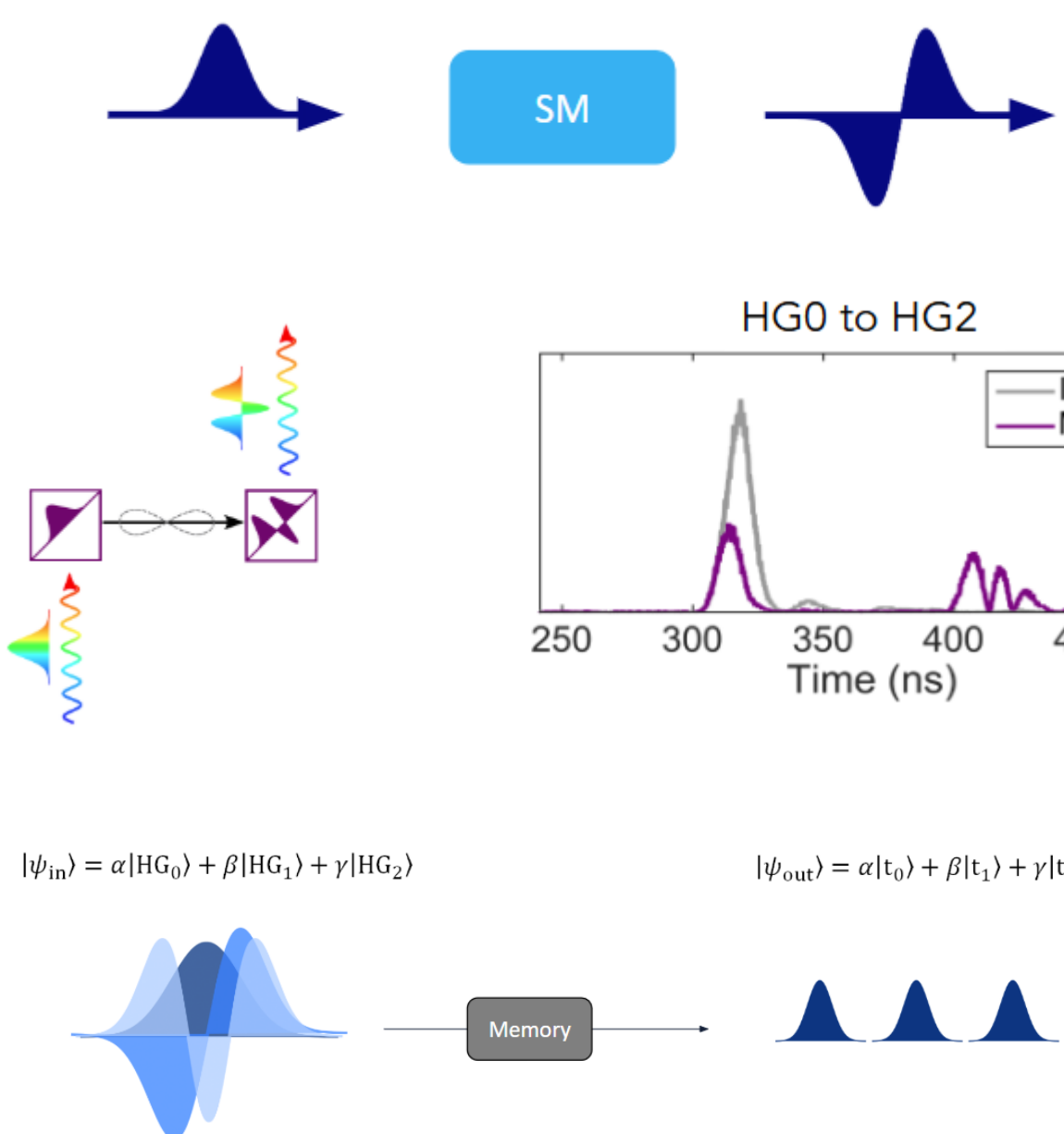
$$|\Theta_n\rangle = \frac{1}{\sqrt{d}} \sum_{m=0}^{d-1} e^{-\frac{2\pi i m n}{d}} |t_m\rangle$$



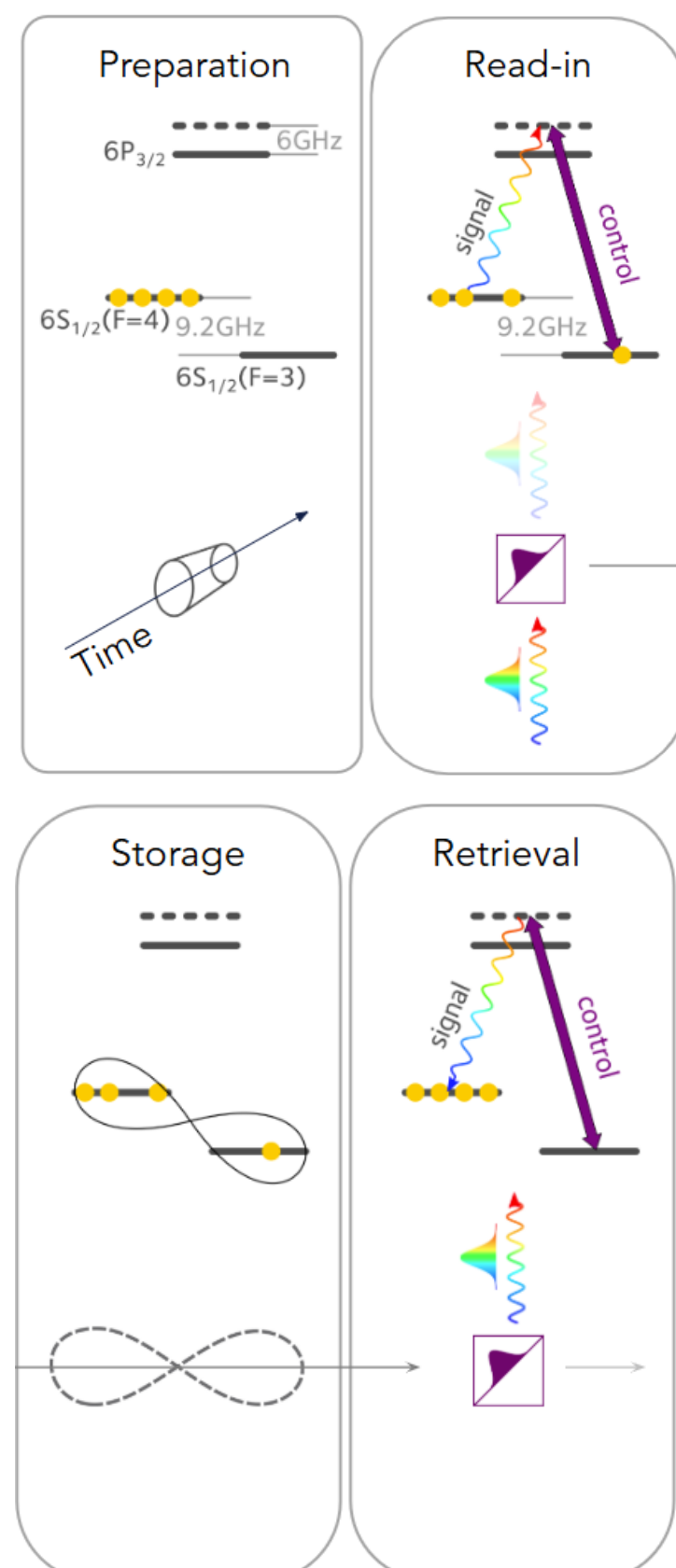
Pulses in time-bins are carved with Mach-Zehnder modulator (MZM) and phases are given by electro-optic phase modulator (EOPM). To detect superpositions pulses are transmitted through group delay dispersion medium (GDD).

A. Widomski, M. Ogrodnik, M. Karpiński, Optica 11, 926 (2024).

Temporal mode manipulation using a quantum memory



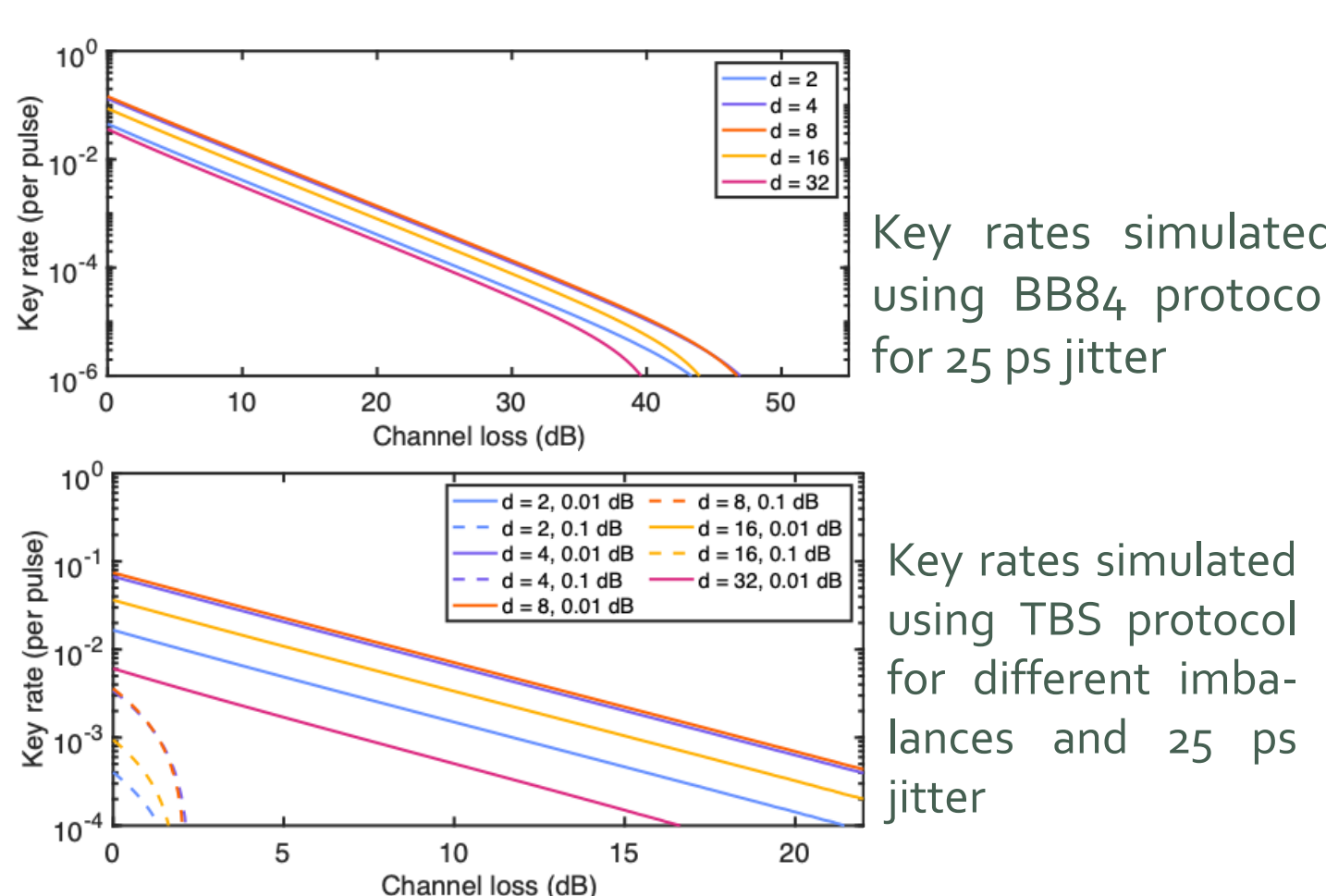
Multipass quantum memory approach was used to convert between overlapping temporal modes and nonoverlapping time bins.



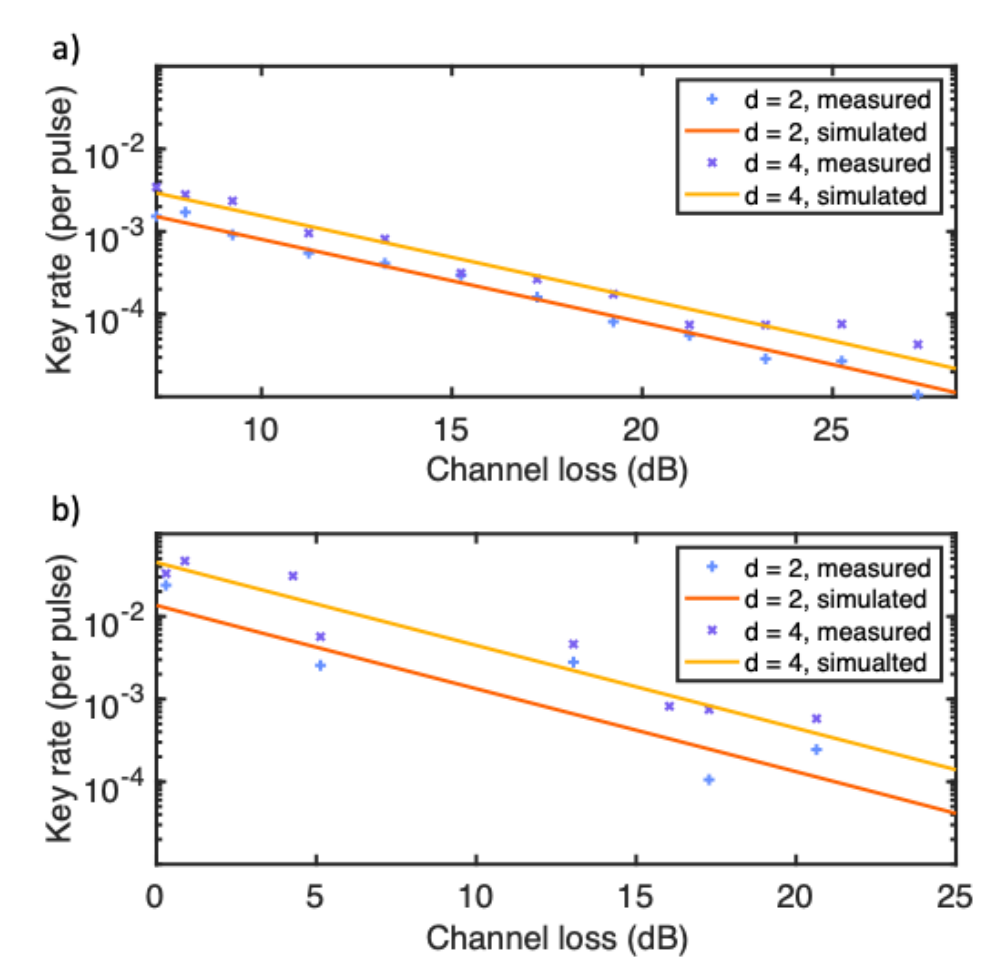
High-dimensional quantum key distribution with resource-efficient detection



By employing a detection scheme based on the temporal Talbot effect, we can perform QKD experiments using only a single detector per measurement, irrespective of the encoding dimension. We present experimentally-obtained secret key rates for both two-dimensional and four-dimensional scenarios, showing that a 4-dimensional encoding offers higher resistance to noise and information entropy than a qubit-based scheme. We show key rates according to two different security proofs – a standard BB84 protocol and a new tunable beam splitter (TBS) protocol [Grasselli2025], developed within QuICHE, which allows dropping assumption about equal detection probability in the two bases.



Ogrodnik et al., Optica Quantum (2025),
Grasselli et al., Phys. Rev. Applied 23, 044011 (2025)



Measured BB84 key rates: a) in-lab
b) over urban dark fibers in Warsaw